

Citation for published version:

Aksulu, MD, Wijers, RAMJ, Van Eerten, H & van der Horst, A 2020, 'A new approach to modelling -ray burst afterglows: Using Gaussian processes to account for the systematics', *Monthly Notices of the Royal Astronomical Society*, vol. 497, no. 4, pp. 4672-4683. <https://doi.org/10.1093/mnras/staa2297>

DOI:

[10.1093/mnras/staa2297](https://doi.org/10.1093/mnras/staa2297)

Publication date:

2020

Document Version

Peer reviewed version

[Link to publication](#)

This is a pre-copyedited, author-produced version of an article accepted for publication in MNRAS following peer review. The version of record is available online at: M D Aksulu, R A M J Wijers, H J van Eerten, A J van der Horst, A new approach to modelling gamma-ray burst afterglows: using Gaussian processes to account for the systematics, *Monthly Notices of the Royal Astronomical Society*, Volume 497, Issue 4, October 2020, Pages 4672–4683, <https://doi.org/10.1093/mnras/staa2297>

University of Bath

Alternative formats

If you require this document in an alternative format, please contact:
openaccess@bath.ac.uk

General rights

Copyright and moral rights for the publications made accessible in the public portal are retained by the authors and/or other copyright owners and it is a condition of accessing publications that users recognise and abide by the legal requirements associated with these rights.

Take down policy

If you believe that this document breaches copyright please contact us providing details, and we will remove access to the work immediately and investigate your claim.

A new approach to modelling gamma-ray burst afterglows: Using Gaussian processes to account for the systematics

M. D. Aksulu,¹★ R. A. M. J. Wijers,¹ H. J. van Eerten,² A. J. van der Horst^{3,4}

¹ Anton Pannekoek Institute for Astronomy, University of Amsterdam, Science Park 904, NL-1098 XH Amsterdam, the Netherlands

² Department of Physics, University of Bath, Claverton Down, Bath BA2 7AY, UK

³ Department of Physics, The George Washington University, 725 21st Street NW, Washington, DC 20052, USA

⁴ Astronomy, Physics, and Statistics Institute of Sciences (APSIS), 725 21st Street NW, Washington, DC 20052, USA

Accepted XXX. Received YYY; in original form ZZZ

ABSTRACT

The afterglow emission from gamma-ray bursts (GRBs) is a valuable source of information to understand the physics of these energetic explosions. The fireball model has become the standard to describe the evolution of the afterglow emission over time and frequency. Thanks to recent developments in the theory of afterglows and numerical simulations of relativistic outflows, we are able to model the afterglow emission with realistic dynamics and radiative processes. Although the models agree with observations remarkably well, the afterglow emission still contains additional physics, instrumental systematics, and propagation effects which make the modelling of these events challenging. In this work, we present a new approach to modelling GRB afterglows, using Gaussian processes (GPs) to take into account systematics in the afterglow data. We show that, using this new approach, it is possible to obtain more reliable estimates of the explosion and microphysical parameters of GRBs. We present fit results for 5 long GRBs and find a preliminary correlation between the isotropic energetics and opening angles of GRBs, which confirms the idea of a common energy reservoir for the kinetic energy of long GRBs.

Key words: gamma-ray burst: general – gamma-ray burst: individual: GRB 970508, 980703, 990510, 991208, 991216 – methods: data analysis – methods: statistical

1 INTRODUCTION

Gamma-ray bursts (GRBs) are the most energetic explosions in the Universe. They are either the result of the collapse of massive stars (long GRBs) (Woosley 1993), or of compact object mergers where at least one of the objects is a neutron star (short GRBs) (Eichler et al. 1989); for a review see, e.g. Piran (2004). During these catastrophic events, an ultra-relativistic, collimated outflow is generated by a compact central engine (Rees & Meszaros 1992). Initially, GRBs are detected as prompt γ -ray flashes. The exact emission mechanism which produces these γ rays is still debated; for a review see, e.g. Kumar & Zhang (2015). As the outflow starts to interact with the circumburst medium, it starts to decelerate and forms a relativistic, collisionless shock where charged particles are accelerated in tangled magnetic fields and emit synchrotron emission across the whole electromagnetic spectrum (Rees & Meszaros 1992). This emission is called the afterglow of the GRB. It is possible to under-

stand more about the physics of GRBs by modelling the afterglow. The afterglow emission reveals how the dynamics of such relativistic shocks evolve over time as well as the microphysical properties in such extreme acceleration regions (Wijers et al. 1997; Sari et al. 1998; Wijers & Galama 1999; Panaitescu & Kumar 2002; Yost et al. 2003).

With the launch of the *Neil Gehrels Swift Observatory* (Gehrels et al. 2004), the detection rate of GRB afterglows has significantly increased. Together with multi-wavelength ground and space based follow-up, and the start of the multi-messenger era, we now have a wealth of data on GRB afterglows (Abbott et al. 2017; MAGIC Collaboration et al. 2019). Moreover, recent advancements in afterglow theory and numerical hydrodynamics allow us to model the dynamics and emission mechanism of GRB afterglows much more reliably (van Eerten 2018). Although the models agree with the general trends of the afterglow data, it is still challenging to get reliable estimates of GRB parameters because of additional physics which is not included in the models (e.g. self-synchrotron Compton scattering effects, reverse shock emission), instrumental systematics, and propagation effects

★ E-mail: m.d.aksulu@uva.nl

(e.g. scintillation in radio, absorption by the host galaxy gas and dust in the optical and X-ray regimes). All these effects introduce systematic deviations to the afterglow observations, and result in a more complex flux evolution over time and frequency than predicted by the models. [Gompertz et al. \(2018\)](#) have shown that there must be intrinsic errors involved when modelling GRB afterglows, by using closure relations to show that the data exhibits inconsistencies which cannot be explained by the measurement errors. In this work, we show that systematic deviations put unrealistically tight constraints on the model parameters when performing parameter estimation where the likelihood function is only proportional to the χ^2 value.

In this paper we introduce a new approach to fitting GRB afterglow data, by modelling the systematics using Gaussian processes (GPs). This way, the model parameters are not bound by artifacts of systematic deviations, and Bayesian parameter estimation gives more reliable parameter uncertainties. In Section 2 we explain the method in detail, in Section 3 we present test results with synthetic data sets, and compare to results obtained by conventional modelling. Moreover, in Section 4 we apply this method to 5 long GRB afterglow data sets and present the results. In Section 5 we further elaborate on the modelling of GRB afterglows, and we conclude in Section 6.

2 METHOD

GPs are a generalization of the Gaussian probability distribution, in the sense that, GPs enable us to define a probability distribution over functions instead of variables or vectors ([Rasmussen & Williams 2006](#)). GPs are non-parametric, stochastic processes, and are therefore a useful tool in regression problems where the underlying model of the data is unknown, as is the case for the systematic differences between GRB afterglow models and observations.

In this section we describe a Gaussian process framework for modelling the systematics in the GRB afterglow data sets. We follow the same methodology described in [Gibson et al. \(2012\)](#), where they used the same approach to model transit light curves of exoplanets, which are affected by significant systematics.

2.1 GRB afterglow data

In order to solve for the many GRB afterglow model parameters, a well-sampled multi-wavelength data set is required. The data set consists of N flux measurements, $\mathbf{y} = (f_{\nu 1}, \dots, f_{\nu N})^T$, measured at times and frequencies $\mathbf{X} = (\mathbf{x}_1, \dots, \mathbf{x}_N)^T = ((t_1, \nu_1)^T, \dots, (t_N, \nu_N)^T)^T$, where \mathbf{X} is an $N \times 2$ matrix. The reported uncertainties of the flux measurements are expressed as $\sigma = (\sigma_1, \dots, \sigma_N)^T$.

2.2 Gaussian process framework

In this work, we use a GP model to take into account any possible systematics in the GRB afterglow data in a non-parametric fashion, where the systematics are described as,

$$f(t, \nu) \sim \mathcal{GP}(\mu(t, \nu, \boldsymbol{\phi}), \Sigma(t, \nu, \boldsymbol{\theta})), \quad (1)$$

where μ is the mean function of the GP (i.e. the afterglow model), Σ is the covariance matrix of the GP model, $\boldsymbol{\phi}$ and $\boldsymbol{\theta}$ represent the GRB parameters and the, so called, *hyperparameters* of the GP model respectively.

The log likelihood of the GP model is described as,

$$\log \mathcal{L}(\mathbf{r}|\mathbf{X}, \boldsymbol{\theta}, \boldsymbol{\phi}) = -\frac{1}{2} \mathbf{r}^T \Sigma^{-1} \mathbf{r} - \frac{1}{2} \log |\Sigma| - \frac{N}{2} \log(2\pi) \quad (2)$$

where \mathbf{r} is the residual of the afterglow model with respect to the observed flux density values. We define the residual as,

$$\mathbf{r} = \log \mathbf{y} - \log \boldsymbol{\mu}, \quad (3)$$

due to the fact that the measured flux densities vary over orders of magnitudes with time and frequency. In such cases it is common to model the logarithm of the measured values ([Snelson et al. 2004](#)). Therefore, we exclude any negative flux measurements from the data sets when modelling.

The covariance matrix, Σ , defines how correlated the data points are over observer time and frequency. In order to construct the covariance matrix, a squared-exponential kernel is chosen, over the 2D input space (time and frequency),

$$\Sigma_{ij} = k(\mathbf{X}_i, \mathbf{X}_j) = A \exp \left[-\frac{1}{2} \sum_{k=1}^2 \frac{(\mathbf{X}_{ik} - \mathbf{X}_{jk})^2}{l_k^2} \right] + \delta_{ij} \sigma_w^2, \quad (4)$$

where A represents the amplitude of the correlations, l_1 and l_2 determine the length scales of the correlations over time and frequency respectively, and σ_w represents the amount of white noise in the data set. These parameters are called the hyperparameters of the GP and need to be marginalized together with the model parameters. The white noise parameter is formulated as $\sigma_w = \sigma_{\log f} \sigma_h$, where $\sigma_{\log f}$ is the uncertainty in the logarithm of the flux measurements and σ_h is the hyperparameter which scales the reported uncertainties. Thus, the hyperparameters can be expressed as,

$$\boldsymbol{\theta} = (A, l_1, l_2, \sigma_h)^T. \quad (5)$$

In this work, we use the **george** Python package ([Ambikasaran et al. 2015](#)) as the GP framework. **george** enables us to calculate efficiently the covariance matrix and the GP likelihood even for relatively large data sets and it is designed to be used with any external optimization/sampling algorithm.

2.3 Model

We assume a relativistic, collimated, outflow interacting with the circumburst medium (CBM), forming a pair of shocks propagating into the ejecta (short-lived reverse shock) and into the CBM (long-lived forward shock) where charged particles are accelerated and emit synchrotron radiation ([Sari et al. 1998](#); [Wijers & Galama 1999](#); [Granot & Sari 2002](#)). In this work we only consider the emission originating from the forward shock. The forward shock model has been able to successfully describe the spectral and temporal evolution of GRB afterglows.

In this work, we incorporate **scalefit** ([Ryan et al. 2015](#), Ryan et al. in prep.), as the mean function of the GP model. **scalefit** is an afterglow model, which makes use of pre-calculated tables of spectral features (i.e self-absorption

break ν_a , injection break ν_m , cooling break ν_c , and peak flux density of the spectrum $f_{\nu, \text{peak}}$ over decades in time and for different observing angles. `scalefit` takes advantage of scale invariance to calculate the observed flux density for given explosion and microphysical parameters, observer times and frequencies. The model parameters are described as,

$$\phi = (\theta_0, E_{K, \text{iso}}, n_0, \theta_{\text{obs}}, p, \epsilon_B, \bar{\epsilon}_e, \xi_N)^T, \quad (6)$$

where θ_0 is the opening angle of the jet, $E_{K, \text{iso}}$ is the isotropic-equivalent kinetic energy of the explosion, n_0 is the circumburst number density, θ_{obs} is the observing angle, p is the power-law index of the accelerated electron population, ϵ_B is the fraction of thermal energy in the magnetic fields, $\bar{\epsilon}_e \equiv \frac{p-2}{p-1} \epsilon_e$ where ϵ_e is the fraction of thermal energy in the accelerated electrons, and ξ_N is the fraction of electrons being accelerated.

`boxfit` (van Eerten et al. 2012) is used to produce the tables containing the spectral features. `boxfit` is a GRB afterglow modelling tool, which makes use of pre-calculated hydrodynamics data and solves radiative transfer equations during runtime. Since it relies on hydrodynamics data, it is able to model the dynamics of the blast wave reliably. `boxfit` has been used to successfully model the broadband emission from various afterglows (see e.g. Guidorzi et al. 2014; Higgs et al. 2019; Kangas et al. 2020), but its computationally expensive repeated radiative transfer calculations are a drawback when implementing `boxfit` in a sampling algorithm. `scalefit` instead draws from a table that reproduces the spectral breaks and peak fluxes from `boxfit` exactly, but approximates the spectral curvature across breaks when reconstructing spectra. Through its approximation of spectral curvature, `scalefit` avoids the need for repeated radiative transfer calculations and allows for fast computation (for applications, see e.g. Ryan et al. 2015; Zhang et al. 2015). This offers a good compromise between speed and accuracy. In this work, aimed at GPs, we make use of the fact that due to the slight differences in their approach the different modelling tools produce afterglow light curves with different relative systematics, leaving a detailed comparison of the relative merit between the methods for afterglow modelling in general for future work.

2.4 Regression

In order to marginalize over the model parameters and the hyperparameters of the GP, we make use of `pymultinest` (Buchner et al. 2014), which is the Python implementation of the MultiNest nested sampling algorithm (Feroz et al. 2009). Sampling from complex objective functions can be challenging as algorithms can get stuck in local maxima. The main advantage of using `pymultinest` is that it is able to converge on the global maximum with high efficiency (i.e. relatively small number of function evaluations). For all the presented results, `pymultinest` is used in the importance sampling mode (Feroz et al. 2019) with mode separation disabled. We use 1000 initial live points and use an evidence tolerance of 0.5 as our convergence criterion. These values are adapted from Feroz et al. (2009).

The fraction of accelerated electrons, ξ_N , is degenerate with respect to $(E_{K, \text{iso}}, n_0, \epsilon_B, \bar{\epsilon}_e)$, where $(E_{K, \text{iso}}, n_0)$ are

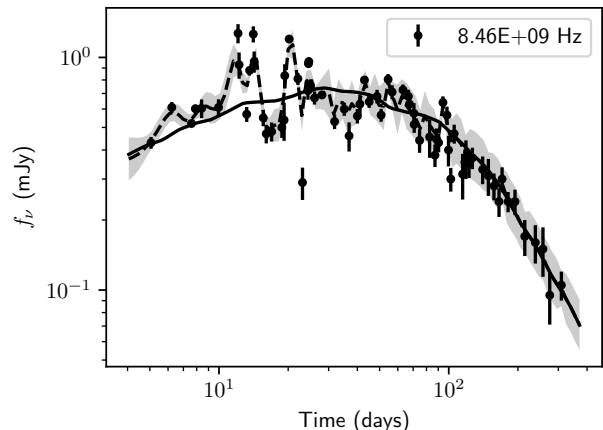


Figure 1. Example regression result from modelling GRB 970508 using the GP model. The radio light curve at 8.46 GHz is shown, where the solid line is the `scalefit` light curve, the dashed line is the mean predicted by the GP model and the shaded area represents the $1\text{-}\sigma$ uncertainty of the GP model. It can be seen that at early times the data is heavily affected by scintillation, and the systematics are modelled by the GP. The variability in the model (solid line) is due to numerical noise.

proportional to $1/\xi_N$, and $(\epsilon_B, \bar{\epsilon}_e)$ are proportional to ξ_N (Eichler & Waxman 2005). Because of this degeneracy, we fix ξ_N to be 0.10. Canonically, ξ_N is set to unity when modelling GRB afterglows, however, for our sample we find that a smaller value for ξ_N gives more physical results for ϵ_B and $\bar{\epsilon}_e$, since accepting the canonical value results in non-physical parameter values (e.g. $\epsilon_B + \epsilon_e > 1$). Moreover, particle-in-cell simulations have shown that ξ_N can be as low as 0.01, depending on the shock conditions (Sironi & Spitkovsky 2011).

Regression is performed by marginalizing over both the hyperparameters and model parameters (see Equations 5 and 6). In all of the fits presented in this work, we assume that the systematics are uncorrelated over the frequency domain by fixing the hyperparameter l_2 to a very small number. We recognize that this assumption may not hold for regions of the spectrum where the frequency domain is sampled closely (e.g. radio or optical observations at similar frequencies). Also, dust extinction may lead to correlated noise in the frequency space over several decades. However, after correcting for dust extinction, when the data set spans over multiple decades in frequency, the emission in radio, optical and X-rays will not be correlated. In Figure 1, we present an example regression result for the radio light curve of GRB 970508, which contains significant variability in radio bands at early times due to interstellar scintillation. It can be seen that over time, as the shock front expands and the source size increases, the variability decreases.

3 APPLICATION TO SYNTHETIC DATA

In order to test the effectiveness of the proposed method, we generate synthetic data sets and try to recover the true parameters by modelling the synthetic data using both the conventional method of sampling the χ^2 likelihood and the

Table 1. Assumed priors for modelling synthetic data sets.

Parameter range	Prior distribution
$0.01 < \theta_0 < 1.6$	log-uniform
$10^{50} < E_{K,\text{iso}} < 10^{55}$	log-uniform
$10^{-4} < n_0 < 1000$	log-uniform
$0 < \theta_{\text{obs}}/\theta_0 < 2$	uniform
$1.5 < p < 3.0$	uniform
$10^{-7} < \epsilon_B < 0.50$	log-uniform
$10^{-4} < \bar{\epsilon}_e < 10$	log-uniform

Table 2. Fit results for the synthetic data set generated using Model 1. The data set contains only white noise as described in Section 3. Results from both χ^2 and GP (\mathcal{GP}) likelihood sampling are presented. All the uncertainties on the parameters represent the 95% credible interval. Parameter estimations which include and exclude the true parameter value within the 95% credible interval are marked as ✓ and ✗, respectively.

Parameter	χ^2	\mathcal{GP}	True value
θ_0	$0.0967^{+0.016}_{-0.0061}$ ✓	$0.098^{+0.017}_{-0.029}$ ✓	0.10
$\log_{10}(E_{K,\text{iso},53})$	$0.01^{+0.17}_{-0.21}$ ✓	$-0.04^{+0.23}_{-0.24}$ ✓	0.00
$\log_{10}(n_0)$	$-0.09^{+0.18}_{-0.23}$ ✓	$-0.12^{+0.26}_{-0.32}$ ✓	0.00
$\theta_{\text{obs}}/\theta_0$	$0.29^{+0.17}_{-0.13}$ ✓	$0.27^{+0.26}_{-0.22}$ ✓	0.30
p	$2.405^{+0.022}_{-0.018}$ ✓	$2.406^{+0.028}_{-0.042}$ ✓	2.40
$\log_{10}(\epsilon_B)$	$-1.94^{+0.27}_{-0.24}$ ✓	$-1.91^{+0.37}_{-0.34}$ ✓	-2.00
$\log_{10}(\bar{\epsilon}_e)$	$-1.48^{+0.15}_{-0.20}$ ✓	$-1.51^{+0.16}_{-0.15}$ ✓	-1.54

proposed method of sampling the GP log likelihood function (Equation 2).

Two sets of synthetic data are generated using **scalefit** (Model 1 from now on) and **boxfit** (Model 2 from now on) as the underlying model. The synthetic data sets are generated in radio, optical and X-ray bands and across 10 time epochs, which are log-uniformly separated. The uncertainty fractions are chosen to be 10%, 2%, and 10%, for radio, optical and X-ray bands respectively. The data points are generated by drawing from a Gaussian distribution with the model value (either Model 1 or 2) as the mean and the corresponding uncertainty as the standard deviation.

In this work, we model any type of data set using Model 1 (see Section 2.3). Therefore, the synthetic data set generated with Model 1 contains only white noise, whereas the synthetic data set generated with Model 2 also contains systematics with respect to Model 1. This allows us to test the performance of the GP model, both in the absence and presence of systematic differences. For all the synthetic data modelling, we use the same prior for the parameters, which is presented in Table 1. We select fiducial GRB parameter values for our synthetic data sets; $(\theta_0, E_{K,\text{iso}}, n_0, \theta_{\text{obs}}/\theta_0, p, \epsilon_B, \epsilon_e) = (0.10, 10^{53}, 1.00, 0.30, 2.4, 10^{-2}, 10^{-1})$.

Tables 2 and 3 show fit results for both data sets and modelling approaches. As it can be seen in Table 2, the GP model and χ^2 sampling perform similarly in the absence of systematics. Overall, the GP model results in larger parameter uncertainties. When there are systematics involved, the

Table 3. Fit results for the synthetic data set generated using Model 2. The data set contains both white noise and systematics as described in Section 3. Results from both χ^2 and GP (\mathcal{GP}) likelihood sampling are presented. All the uncertainties on the parameters represent the 95% credible interval. Parameter estimations which include and exclude the true parameter value within the 95% credible interval are marked as ✓ and ✗, respectively.

Parameter	χ^2	\mathcal{GP}	True value
θ_0	$0.03570^{+0.00081}_{-0.00072}$ ✗	$0.071^{+0.062}_{-0.039}$ ✓	0.10
$\log_{10}(E_{K,\text{iso},53})$	$0.519^{+0.034}_{-0.034}$ ✗	$0.57^{+0.87}_{-1.0}$ ✓	0.00
$\log_{10}(n_0)$	$-2.669^{+0.044}_{-0.048}$ ✗	$-0.2^{+1.2}_{-1.3}$ ✓	0.00
$\theta_{\text{obs}}/\theta_0$	$0.641^{+0.033}_{-0.038}$ ✗	$0.70^{+0.46}_{-0.58}$ ✓	0.30
p	$2.469^{+0.014}_{-0.014}$ ✗	$2.33^{+0.11}_{-0.12}$ ✓	2.40
$\log_{10}(\epsilon_B)$	$-0.538^{+0.081}_{-0.081}$ ✗	$-2.0^{+1.5}_{-1.6}$ ✓	-2.00
$\log_{10}(\bar{\epsilon}_e)$	$-1.779^{+0.018}_{-0.017}$ ✗	$-1.96^{+0.46}_{-0.52}$ ✓	-1.54

shortcomings of the χ^2 sampling approach stand out. In Table 3, we show that the χ^2 sampling technique is unable to recover any of the true parameters, despite inferring small uncertainties on the parameters. On the other hand the GP model is able to recover every parameter within the 95% credible interval.

In Figure 2 we show the fit results for all synthetic data sets and modelling approaches in the form of violin plots. Violin plots are a way to visualize the marginalized distributions of parameters in a compact way, where the shaded area represents the normalized histogram of the posterior samples, and the solid bar shows the interquartile range of the distribution. As it can be seen, the GP model results in larger uncertainties, more complex marginal distributions, and more accurate parameter estimations.

In order to investigate further whether sampling the GP likelihood results in more reliable parameter inferences, regardless of the chosen parameters for the synthetic data sets, we generate 100 sets of synthetic data both using Model 1 and Model 2 with randomly chosen GRB parameters. We fit all of the data sets using both the proposed method of GP likelihood sampling and χ^2 sampling. We perform coverage measurements on these fit results to determine how accurate the inferred uncertainties are. Coverage measurements are performed by fitting 100 synthetic data sets and counting how many times the true parameter was recovered for a given confidence region.

In Figure 3, we show the coverage measurement results for each GRB parameter both in the presence and absence of systematic deviations. These plots show the fraction of successfully recovered parameters (vertical axis) for a given credible interval ¹ (horizontal axis) (Sellentin & Starck 2019). The black points show the ideally expected coverage, where the error bars are calculated using the binomial uncertainty, given by,

$$\sigma = \sqrt{p(1-p)/N} \quad (7)$$

where p is the probability of containing the true parameter (credible interval) and N is the number of samples (100 in

¹ Credible interval is the Bayesian analogue of confidence interval.

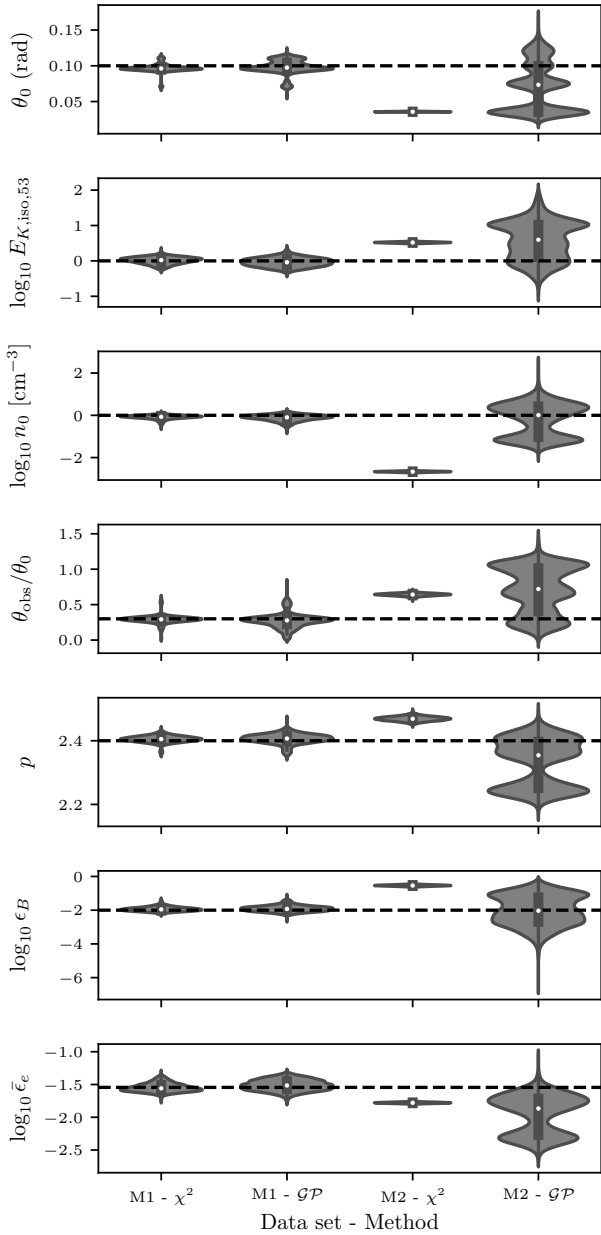


Figure 2. Violin plot of the fit results for different synthetic data sets and methods. Data sets which are generated by Model 2 (M2) contain systematics and white noise. Data sets generated with Model 1 (M1) contain no systematics but the same amount of white noise as Model 2. Both data sets are fitted using χ^2 sampling (χ^2) and sampling the GP likelihood in Equation 2 (\mathcal{GP}). The horizontal dashed lines represent the true parameter values.

our case). The coverage measurements show that the GP model performs better both in the presence and absence of systematic deviations, as the measured coverage for the GP model is closer to the ideal case. χ^2 sampling underestimates the errors on the parameters, especially for parameters which affect the temporal slope of the light curves (θ_0 , θ_{obs} and p). In the presence of systematic deviations, even

Table 4. Redshift (z), Galactic foreground extinction ($A_{V,\text{MW}}$), rest-frame host galaxy extinction ($A_{V,\text{host}}$), and the best fit extinction model for the host galaxy of the long GRB sample. MW denotes Milky Way type host extinction, and SMC denotes Small Magellanic Cloud type host extinction (Pei 1992).

GRB name	z	$A_{V,\text{MW}}$	$A_{V,\text{host}}$	Host type
970508	0.835	~ 0	~ 0	N/A
980703	0.966	0.1891	0.90	MW
990510	1.619	~ 0	0.22	SMC
991208	0.706	0.0512	0.80	MW
991216	1.02	2.016	~ 0	N/A

Table 5. Assumed priors for modelling the long GRB sample.

Parameter range	Prior distribution
$0.01 < \theta_0 < 1.6$	log-uniform
$10^{48} < E_{K,\text{iso}} < 10^{55}$	log-uniform
$10^{-4} < n_0 < 1000$	log-uniform
$0 < \theta_{\text{obs}}/\theta_0 < 2$	uniform
$1.0 < p < 3.0$	uniform
$10^{-7} < \epsilon_B < 0.50$	log-uniform
$10^{-5} < \epsilon_e < 10$	log-uniform

the GP model underestimates the errors on the parameters, however, less so than χ^2 sampling.

4 APPLICATION TO ARCHIVAL GRB AFTERGLOW DATA

In this section we present fit results for 5 long GRB afterglows for which significant modeling has already been done, namely; GRB 970508, GRB 980703, GRB 990510, GRB 991208 and GRB 991216. We compare our results to previous, multi-wavelength, modelling efforts. In Table 4, we present the overall properties of the GRB sample. For this work, we were able to generate `scalefit` tables for constant density CBM only. Therefore, constant density CBM is assumed when fitting the afterglow data. We use the prior distributions presented in Table 5 in our modelling efforts. The inferred parameter distributions for the long GRB sample can be seen in Figure 4 in the form of a violin plot.

4.1 GRB 970508

GRB 970508 exhibits an increase in optical flux at around ~ 1 day after the burst, and starts to decline as a power-law with time. Panaitescu & Kumar (2002) (PK02 from now on) explain the rise in the optical flux by assuming that the jet is viewed off-axis with $\theta_{\text{obs}} \sim 4/3\theta_0$. They find that a wind-like CBM ($n \propto r^{-2}$) suits the observations best. On the other hand, Yost et al. (2003) (Y03 from now on) favor a constant density CBM in their analysis.

In our analysis, we exclude the data points before the peak of the rise in optical wavelengths, and fit the data points which obey the power-law behaviour. We take the

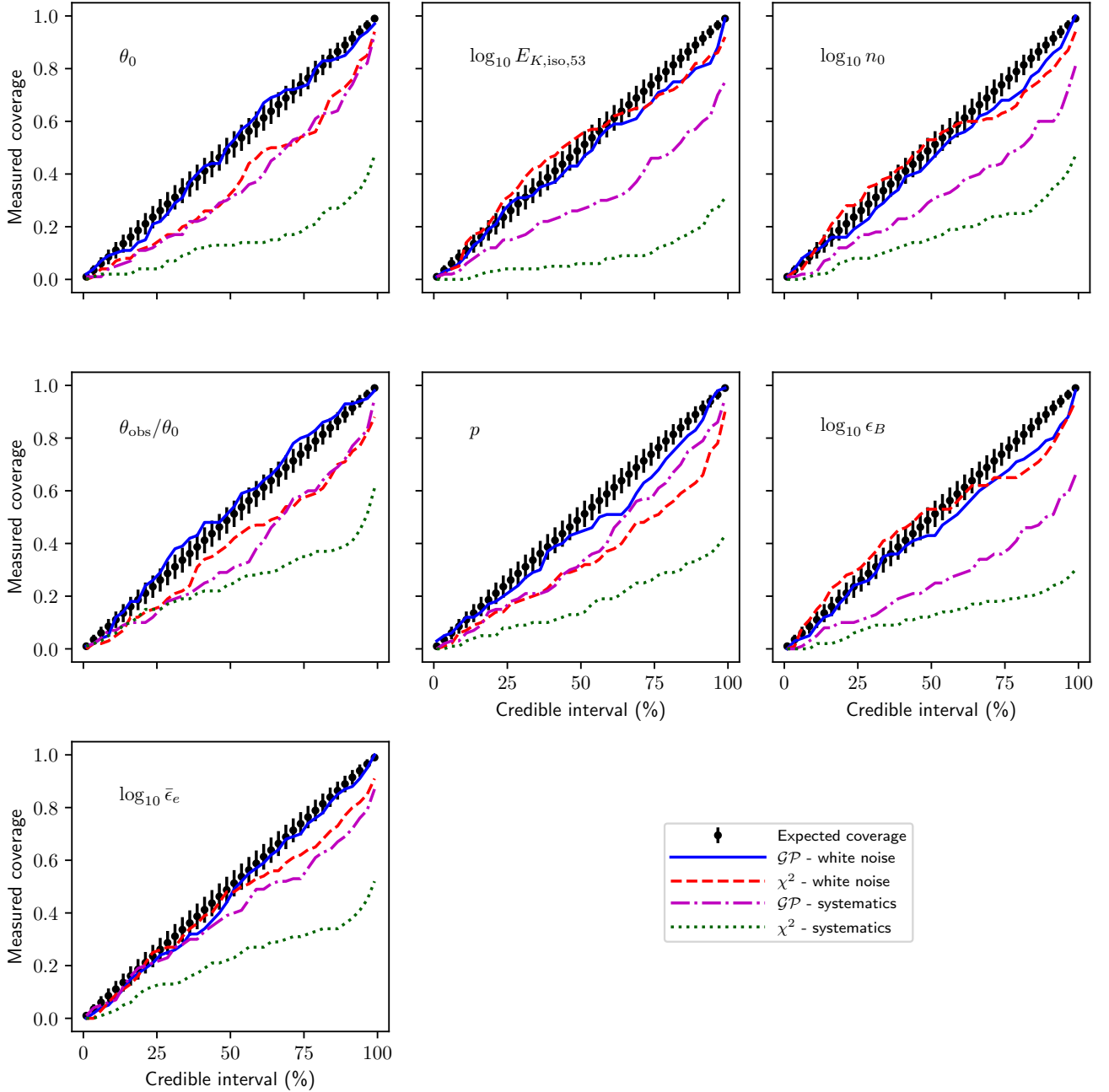


Figure 3. Coverage measurement results for GRB parameters. The blue and red lines show the coverage measurement results in the case where the data only contains white noise (synthetic data set generated by Model 1) for GP model regression (\mathcal{GP}) and χ^2 sampling, respectively. The magenta and yellow lines show the coverage measurement results in the case where the data contains systematic deviations (synthetic data set generated by Model 2) for GP model regression and χ^2 sampling, respectively. The black points show the ideally expected coverage, where the error bars are calculated using the binomial uncertainty $\sigma = \sqrt{p(1-p)/N}$, where p is the probability of containing the true parameter (credible interval) and N is the number of samples (100 in our case).

observing angle as a free parameter, allowing viewing angles both larger and smaller than the opening angle.

Figure A1 shows the light curves for the inferred parameter distribution for the afterglow of GRB 970508 using GP regression. Table 6 shows the inferred parameter values. Assuming that the re-brightening in optical bands is due to late-time energy injection from the central engine, the ener-

getics inferred from our modelling will overestimate the initial explosion energy. We find a wide opening angle, $0.74^{+0.52}_{-0.28}$ rad, which is consistent with Y03, who find an opening angle of $0.84^{+0.03}_{-0.03}$ rad. The inferred isotropic kinetic energy by the GP model is consistent with both PK02 and Y03, whereas the χ^2 sampling infers a lower $E_{K,\text{iso}}$ with small uncertainty.

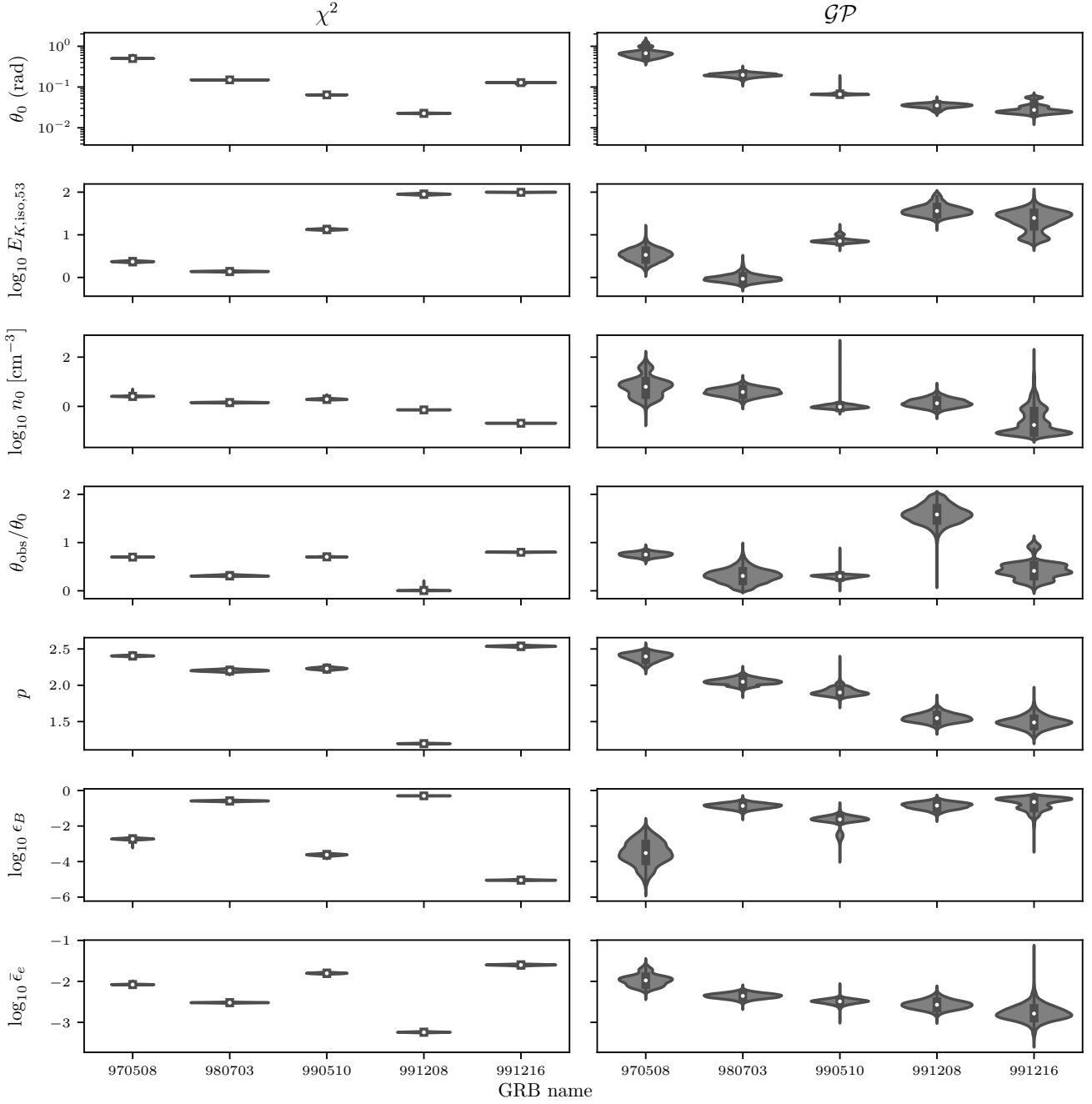


Figure 4. Violin plot showing the inferred marginalized distributions of the GRB parameters for the long GRB sample. The panel on the left shows the results for χ^2 sampling whereas the panel on the right shows the results from GP likelihood sampling.

Both GP likelihood and χ^2 sampling infer similar p values of ~ 2.4 , which is larger than what PK02 and Y03 found.

4.2 GRB 980703

Panaitescu & Kumar (2001) (PK01 from now on) favor a constant density CBM for GRB 980703. Following Vreeswijk et al. (1999), PK01 take the host extinction to be $A_V = 1.45 \pm 0.13$, and find that $p = 3.08$. Y03 also favor a constant density CBM, and find a value of $A_V = 1.15$ for the host

extinction and infer $p = 2.54^{+0.04}_{-0.1}$. In this work we take the host extinction to be $A_V = 0.9$ (Bloom et al. 1998) and find that $p = 2.05^{+0.10}_{-0.095}$.

The host galaxy of GRB 980703 has a significant contribution to the observed radio and optical emission. We assume that the host galaxy contribution is constant over time and leave the host galaxy flux in radio and optical wavelengths as free parameters.

Figure A2 shows the light curves for the inferred parameter distribution for the afterglow of GRB 980703 using GP regression. Table 7 shows the inferred parameter values.

Table 6. Fit results for GRB 970508. Results from both χ^2 sampling and GP likelihood sampling (\mathcal{GP}) are presented. The uncertainties on the parameters represent the 95% credible interval for columns χ^2 and \mathcal{GP} . Columns PK02 and Y03 show results from Panaitescu & Kumar (2002) and Yost et al. (2003), respectively. PK02 did not provide uncertainties on the parameters for this burst. The uncertainties for the Y03 results represent the 68.3% credible interval. All the values taken from previous studies have been converted to the same units. The values have been corrected for our choice of $\xi_N = 0.1$ by multiplying ($E_{K,\text{iso}}$, n_0) by 10 and dividing (ϵ_B , $\bar{\epsilon}_e$) by 10 (see Section 2.4).

GRB 970508				
Parameter	χ^2	\mathcal{GP}	PK02	Y03
θ_0	$0.5007^{+0.0063}_{-0.0055}$	$0.74^{+0.52}_{-0.28}$	0.32	$0.84^{+0.030}_{-0.030}$
$\log_{10}(E_{K,\text{iso},53})$	$0.370^{+0.027}_{-0.026}$	$0.53^{+0.28}_{-0.28}$	0.597	$0.56^{+0.011}_{-0.011}$
$\log_{10}(n_0)$	$0.407^{+0.028}_{-0.027}$	$0.80^{+0.90}_{-0.69}$	0.87	$0.30^{+0.021}_{-0.045}$
$\theta_{\text{obs}}/\theta_0$	$0.6977^{+0.0055}_{-0.0074}$	$0.750^{+0.093}_{-0.096}$	1.33	0
p	$2.404^{+0.013}_{-0.013}$	$2.39^{+0.10}_{-0.12}$	2.18	$2.12^{+0.03}_{-0.008}$
$\log_{10}(\epsilon_B)$	$-2.729^{+0.075}_{-0.077}$	$-3.5^{+1.2}_{-1.4}$	-2.34	$-1.60^{+0.010}_{-0.036}$
$\log_{10}(\bar{\epsilon}_e)$	$-2.078^{+0.011}_{-0.011}$	$-1.97^{+0.30}_{-0.25}$	-2.77	$-2.43^{+0.02}_{-0.015}$

Table 7. Fit results for GRB 980703. Results from both χ^2 sampling and GP likelihood sampling (\mathcal{GP}) are presented. See Table 6 for detailed explanation.

GRB 980703			
Parameter	χ^2	\mathcal{GP}	Y03
θ_0	$0.14990^{+0.00093}_{-0.0011}$	$0.199^{+0.043}_{-0.042}$	$0.234^{+0.02}_{-0.007}$
$\log_{10}(E_{K,\text{iso},53})$	$0.139^{+0.016}_{-0.015}$	$-0.02^{+0.19}_{-0.17}$	$1.07^{+0.028}_{-0.080}$
$\log_{10}(n_0)$	$0.156^{+0.026}_{-0.025}$	$0.58^{+0.33}_{-0.34}$	$2.44^{+0.057}_{-0.049}$
$\theta_{\text{obs}}/\theta_0$	$0.313^{+0.023}_{-0.018}$	$0.31^{+0.28}_{-0.29}$	0
p	$2.202^{+0.026}_{-0.026}$	$2.049^{+0.092}_{-0.094}$	$2.54^{+0.04}_{-0.1}$
$\log_{10}(\epsilon_B)$	$-0.588^{+0.048}_{-0.049}$	$-0.87^{+0.28}_{-0.31}$	$-3.74^{+0.087}_{-0.079}$
$\log_{10}(\bar{\epsilon}_e)$	$-2.520^{+0.015}_{-0.015}$	$-2.36^{+0.13}_{-0.13}$	$-2.02^{+0.065}_{-0.110}$

The GP model infers an opening angle which is consistent with Y03, whereas χ^2 likelihood sampling infers a smaller opening angle. Inferred n_0 and $E_{K,\text{iso}}$ values are significantly smaller than Y03.

4.3 GRB 990510

PK01 favor a constant density CBM for the case of GRB 990510. The optical afterglow of GRB 990510 exhibits a break in its temporal evolution at around 1.5 days. This break is interpreted as a jet-break. PK01 find $p = 2.09 \pm 0.03$ using closure relations, which is also consistent with the inferred value from the GP model.

Figure A3 shows the light curves for the inferred parameter distribution for the afterglow of GRB 990510 using GP regression. Table 8 shows the inferred parameter values. The inferred opening angle is consistent with van Eerten et al. (2012), where they performed a detailed fit using `box-fit` (vE12) model and found an opening angle $0.075^{+0.002}_{-0.004}$ rad assuming an on-axis observer. On the other hand PK02 find a smaller opening angle $0.054^{+0.001}_{-0.006}$ rad. The GP model predicts a larger ϵ_B value when compared to χ^2 sampling and previous studies.

4.4 GRB 991208

Figure A4 shows the light curves for the inferred parameter distribution for the afterglow of GRB 991208 using GP regression. Table 9 shows the inferred parameter values. The inferred p value by the GP model agrees with the results presented in PK02, whereas χ^2 sampling results in a smaller p value.

In our analysis we find a smaller opening angle than PK02 with an extremely off-axis observer angle. The GP regression and χ^2 sampling give significantly different results for microphysical parameters and observer angle.

4.5 GRB 991216

Figure A5 shows the light curves for the inferred parameter distribution for the afterglow of GRB 991216 using GP regression. Table 10 shows the inferred parameter values. PK02 find a hard electron distribution with $p = 1.36 \pm 0.03$, which is consistent with the results we get from GP modelling.

GP regression and χ^2 sampling result in very different parameter values for GRB 991216. This is mainly because the optical data contribute to the χ^2 value the most, whereas the radio data has a small contribution to the χ^2 . The best fit obtained from χ^2 sampling, despite the fact that it has a smaller χ^2 value than the best fit of the GP model, completely misses the radio data points and therefore is an inadequate representation of the observed emission.

5 DISCUSSION

GRBs are thought to be collimated outflows, therefore the isotropic equivalent energies of these events are an overestimation of the true energetics. The true, beaming corrected, energies of these events significantly depend on the geometry of the outflow (i.e. the opening angle),

$$E_K = E_{K,\text{iso}}(1 - \cos \theta_0). \quad (8)$$

Table 8. Fit results for GRB 990510. Results from both χ^2 sampling and GP likelihood sampling (\mathcal{GP}) are presented. The uncertainties on the parameters represent the 95% credible interval for columns χ^2 and \mathcal{GP} . Columns PK02 and vE12 show results from Panaitescu & Kumar (2002) and van Eerten et al. (2012). The uncertainties for the PK02 and vE12 results represent the 90% and 68% credible intervals, respectively. All the values taken from previous studies have been converted to the same units. The values have been corrected for our choice of $\xi_N = 0.1$ by multiplying $(E_{K,\text{iso}}, n_0)$ by 10 and dividing $(\epsilon_B, \bar{\epsilon}_e)$ by 10 (see Section 2.4). Since PK02 find $p < 2$, the conversion from ϵ_e to the equivalent $\bar{\epsilon}_e$ results in a negative value. Therefore we denote $\log_{10}(\bar{\epsilon}_e)$ as N/A.

GRB 990510				
Parameter	χ^2	\mathcal{GP}	PK02	vE12
θ_0	$0.06423^{+0.00078}_{-0.00077}$	$0.0671^{+0.0062}_{-0.0042}$	$0.054^{+0.0017}_{-0.0087}$	$0.075^{+0.002}_{-0.004}$
$\log_{10}(E_{K,\text{iso},53})$	$1.125^{+0.026}_{-0.024}$	$0.870^{+0.18}_{-0.089}$	$0.98^{+0.80}_{-0.21}$	$1.25^{+0.06}_{-0.02}$
$\log_{10}(n_0)$	$0.293^{+0.049}_{-0.042}$	$-0.01^{+0.15}_{-0.13}$	$0.46^{+0.139}_{-0.316}$	$-0.52^{+0.054}_{-0.221}$
$\theta_{\text{obs}}/\theta_0$	$0.705^{+0.014}_{-0.0093}$	$0.297^{+0.057}_{-0.074}$	0	0
p	$2.230^{+0.028}_{-0.027}$	$1.91^{+0.12}_{-0.089}$	$1.83^{+0.18}_{-0.01}$	$2.28^{+0.06}_{-0.01}$
$\log_{10}(\epsilon_B)$	$-3.620^{+0.086}_{-0.089}$	$-1.72^{+0.42}_{-1.0}$	$-3.28^{+0.95}_{-1.01}$	$-3.33^{+0.07}_{-0.08}$
$\log_{10}(\bar{\epsilon}_e)$	$-1.803^{+0.026}_{-0.026}$	$-2.49^{+0.11}_{-0.12}$	N/A	$-2.08^{+0.13}_{-0.02}$

Table 9. Fit results for GRB 991208. Results from both χ^2 sampling and GP likelihood sampling (\mathcal{GP}) are presented together with literature values. See Table 8 for detailed explanation.

GRB 991208			
Parameter	χ^2	\mathcal{GP}	PK02
θ_0	$0.02261^{+0.00059}_{-0.00059}$	$0.0350^{+0.0072}_{-0.0088}$	$0.22^{+0.026}_{-0.038}$
$\log_{10}(E_{K,\text{iso},53})$	$1.951^{+0.024}_{-0.024}$	$1.58^{+0.30}_{-0.23}$	$-0.015^{+0.49}_{-0.27}$
$\log_{10}(n_0)$	$-0.1410^{+0.0073}_{-0.0077}$	$0.14^{+0.36}_{-0.32}$	$2.25^{+0.34}_{-0.17}$
$\theta_{\text{obs}}/\theta_0$	$0.0079^{+0.017}_{-0.0083}$	$1.58^{+0.38}_{-0.34}$	0
p	$1.1964^{+0.0077}_{-0.0074}$	$1.55^{+0.13}_{-0.12}$	$1.53^{+0.03}_{-0.03}$
$\log_{10}(\epsilon_B)$	$-0.30186^{+0.00086}_{-0.0018}$	$-0.86^{+0.35}_{-0.38}$	$-2.45^{+0.43}_{-0.39}$
$\log_{10}(\bar{\epsilon}_e)$	$-3.241^{+0.017}_{-0.016}$	$-2.57^{+0.23}_{-0.21}$	N/A

Table 10. Fit results for GRB 991216. Results from both χ^2 sampling and GP likelihood sampling (\mathcal{GP}) are presented together with literature values. See Table 8 for detailed explanation.

GRB 991216			
Parameter	χ^2	\mathcal{GP}	PK02
θ_0	$0.1268^{+0.0031}_{-0.0095}$	$0.033^{+0.028}_{-0.012}$	$0.047^{+0.006}_{-0.017}$
$\log_{10}(E_{K,\text{iso},53})$	$1.9961^{+0.0041}_{-0.0082}$	$1.34^{+0.36}_{-0.51}$	$0.99^{+0.68}_{-0.31}$
$\log_{10}(n_0)$	$-0.6833^{+0.0094}_{-0.010}$	$-0.60^{+1.0}_{-0.64}$	$1.67^{+0.38}_{-0.20}$
$\theta_{\text{obs}}/\theta_0$	$0.7979^{+0.0071}_{-0.0095}$	$0.43^{+0.54}_{-0.33}$	0
p	$2.536^{+0.016}_{-0.017}$	$1.49^{+0.18}_{-0.16}$	$1.36^{+0.03}_{-0.03}$
$\log_{10}(\epsilon_B)$	$-5.041^{+0.044}_{-0.036}$	$-0.76^{+0.46}_{-0.70}$	$-2.74^{+0.46}_{-0.21}$
$\log_{10}(\bar{\epsilon}_e)$	$-1.601^{+0.025}_{-0.028}$	$-2.77^{+0.42}_{-0.36}$	N/A

Previous studies have shown that there is observational evidence that there exists a standard energy reservoir for GRBs. Frail et al. (2001) have measured the opening angle of a sample of GRBs based on achromatic breaks in the afterglow light curve. They have shown that the beaming corrected energy release in γ -rays is narrowly clustered around 5×10^{50}

erg. Moreover, Panaitescu & Kumar (2002) have shown, using multi-wavelength afterglow modelling, that the beaming corrected kinetic energy of GRBs are narrowly distributed and vary between 10^{50} to 5×10^{50} erg. Similarly, Berger et al. (2003) have analysed the X-ray afterglow data for a large sample of GRBs with known jet breaks, and have found evidence that the beaming corrected kinetic energy of these events are approximately constant. Note that, these studies are based on pre-*Swift* afterglow observations and therefore might be biased towards more energetic GRBs. After the launch of *Swift*, thanks to improved localization, the number of redshift measurements have increased and other classes of GRBs have been discovered, such as low-luminosity GRBs (ll-GRBs) which exhibit significantly less luminous prompt emission. ll-GRBs might constitute outliers in the overall GRB population and might not conform with the idea of a constant energy reservoir (Liang et al. 2007).

Our analysis also shows a strong correlation between θ_0 and $E_{K,\text{iso}}$. In Figure 5, we show the measured opening angles and isotropic energies of our GRB sample. It can be seen that, when the GP model is used for inferring parameters, the measured values suggest that the beaming corrected kinetic energies are approximately the same for long GRBs. GRB 970508 is a clear outlier, which is consistent with the findings of Panaitescu & Kumar (2002). As discussed in Section 4.1, GRB 970508 exhibits a re-brightening in optical wavelengths, which could be due to late-time energy injection. This energy injection could account for the overestimation of the isotropic-equivalent kinetic energy of this source, which might imply that the standard energy reservoir applies more strongly to the initial ejecta formation of a GRB than to any later activity of the central engine. Note that the correlation is not apparent when χ^2 sampling is used for parameter estimation.

In Figure 6, we show the inferred beaming corrected kinetic energies with 95% credible intervals for the sample GRBs. The inferred E_K in our analysis is $\sim 1.7 \times 10^{51}$ erg, which is about an order of magnitude larger than what previous studies have found. This discrepancy with previous studies is expected as we fix ξ_N to be 0.1 instead of the canonical value of 1.0. We also recognise that it is too early

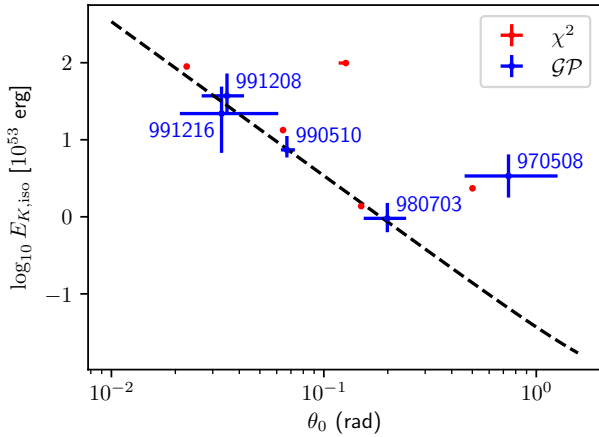


Figure 5. Isotropic equivalent kinetic energy ($E_{K,\text{iso}}$) dependence on the opening angle (θ_0) inferred from our modelling. The red measurements (χ^2) are obtained by χ^2 sampling and the blue measurements (\mathcal{GP}) are obtained by sampling the GP log likelihood function. The dashed black line represents the $E_{K,\text{iso}}(1 - \cos \theta_0) = 1.7 \times 10^{51}$ relation. The error bars represent the 95% credible interval.

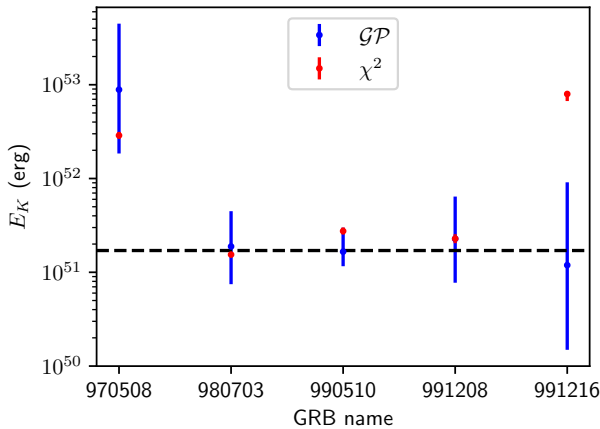


Figure 6. Beaming corrected kinetic energies of the long GRB sample. The red measurements (χ^2) are obtained by χ^2 sampling and the blue measurements (\mathcal{GP}) are obtained by sampling the GP log likelihood function. The dashed line is the log-average of GRBs 980703, 990510, 991208, 991216, which is equal to 1.7×10^{51} erg. The error bars represent the 95% credible interval.

to judge whether these few very well studied, well-sampled GRB afterglows are representative of the whole population.

6 CONCLUSION

In this work, we have introduced a novel method for modelling GRB afterglows, where Gaussian processes are used to take into account any systematics between the model and observations in a non-parametric fashion. Using synthetic data sets, we have shown that the GP approach results in

more accurate posterior distributions with respect to sampling the χ^2 likelihood.

We model a sample of 5 well-known long GRBs with multi-wavelength coverage (GRBs 970508, 980703, 990510, 991208, 991216), using the `scalefit` code together with the GP framework. We compare the inferred parameters for each GRB with the literature values and comment upon the parameter distributions of the overall sample. We find a correlation between the isotropic-kinetic energy and opening angle, with GRB 970508 being the only outlier. This correlation, which is consistent with previous studies, suggests that there is a common energy reservoir which drives the dynamics of GRBs.

ACKNOWLEDGEMENTS

We would like to thank Vatsal Panwar for valuable discussions on Gaussian processes. We would like to thank Dr. Elena Sellentin for her useful tips regarding coverage measurements. We would also like to thank the anonymous referee for their detailed comments. H. J. van Eerten acknowledges partial support by the European Union Horizon 2020 Programme under the AHEAD2020 project (grant agreement number 871158). This work was sponsored by NWO Exact and Natural Sciences for the use of supercomputer facilities.

DATA AVAILABILITY

The data underlying this article can be shared on reasonable request to the corresponding author.

REFERENCES

- Abbott B. P., et al., 2017, *Physical Review Letters*, **119**, 161101
- Ambikasaran S., Foreman-Mackey D., Greengard L., Hogg D. W., O’Neil M., 2015, *IEEE Transactions on Pattern Analysis and Machine Intelligence*, **38**, 252
- Berger E., Kulkarni S. R., Frail D. A., 2003, *ApJ*, **590**, 379
- Bloom J. S., et al., 1998, *ApJ*, **508**, L21
- Buchner J., et al., 2014, *A&A*, **564**, A125
- Eichler D., Waxman E., 2005, *ApJ*, **627**, 861
- Eichler D., Livio M., Piran T., Schramm D. N., 1989, *Nature*, **340**, 126
- Feroz F., Hobson M. P., Bridges M., 2009, *MNRAS*, **398**, 1601
- Feroz F., Hobson M. P., Cameron E., Pettitt A. N., 2019, *The Open Journal of Astrophysics*, **2**, 10
- Frail D. A., et al., 2001, *ApJ*, **562**, L55
- Gehrels N., et al., 2004, *ApJ*, **611**, 1005
- Gibson N. P., Aigrain S., Roberts S., Evans T. M., Osborne M., Pont F., 2012, *MNRAS*, **419**, 2683
- Gompertz B. P., Fruchter A. S., Pe’er A., 2018, *ApJ*, **866**, 162
- Granot J., Sari R., 2002, *ApJ*, **568**, 820
- Guidorzi C., et al., 2014, *MNRAS*, **438**, 752
- Higgins A. B., et al., 2019, *MNRAS*, **484**, 5245
- Kangas T., et al., 2020, *ApJ*, **894**, 43
- Kumar P., Zhang B., 2015, *Phys. Rep.*, **561**, 1
- Liang E., Zhang B., Virgili F., Dai Z. G., 2007, *ApJ*, **662**, 1111
- MAGIC Collaboration et al., 2019, *Nature*, **575**, 455
- Panaiteanu A., Kumar P., 2001, *ApJ*, **554**, 667
- Panaiteanu A., Kumar P., 2002, *ApJ*, **571**, 779
- Pei Y. C., 1992, *ApJ*, **395**, 130

- Piran T., 2004, *Reviews of Modern Physics*, **76**, 1143
- Rasmussen C. E., Williams C. K. I., 2006, *Gaussian Processes for Machine Learning*
- Rees M. J., Meszaros P., 1992, *MNRAS*, **258**, 41
- Ryan G., van Eerten H., MacFadyen A., Zhang B.-B., 2015, *ApJ*, **799**, 3
- Sari R., Piran T., Narayan R., 1998, *ApJ*, **497**, L17
- Sellentin E., Starck J.-L., 2019, *J. Cosmology Astropart. Phys.*, **2019**, 021
- Sironi L., Spitkovsky A., 2011, *ApJ*, **726**, 75
- Snelson E., Ghahramani Z., Rasmussen C. E., 2004, in Thrun S., Saul L. K., Schölkopf B., eds, *Advances in Neural Information Processing Systems 16*. MIT Press, pp 337–344, <http://papers.nips.cc/paper/2481-warped-gaussian-processes.pdf>
- Vreeswijk P. M., et al., 1999, *ApJ*, **523**, 171
- Wijers R. A. M. J., Galama T. J., 1999, *ApJ*, **523**, 177
- Wijers R. A. M. J., Rees M. J., Meszaros P., 1997, *MNRAS*, **288**, L51
- Woosley S. E., 1993, *ApJ*, **405**, 273
- Yost S. A., Harrison F. A., Sari R., Frail D. A., 2003, *ApJ*, **597**, 459
- Zhang B.-B., van Eerten H., Burrows D. N., Ryan G. S., Evans P. A., Racusin J. L., Troja E., MacFadyen A., 2015, *ApJ*, **806**, 15
- van Eerten H., 2018, *International Journal of Modern Physics D*, **27**, 1842002
- van Eerten H., van der Horst A., MacFadyen A., 2012, *ApJ*, **749**, 44

APPENDIX A: FIT RESULTS

This paper has been typeset from a \LaTeX file prepared by the author.

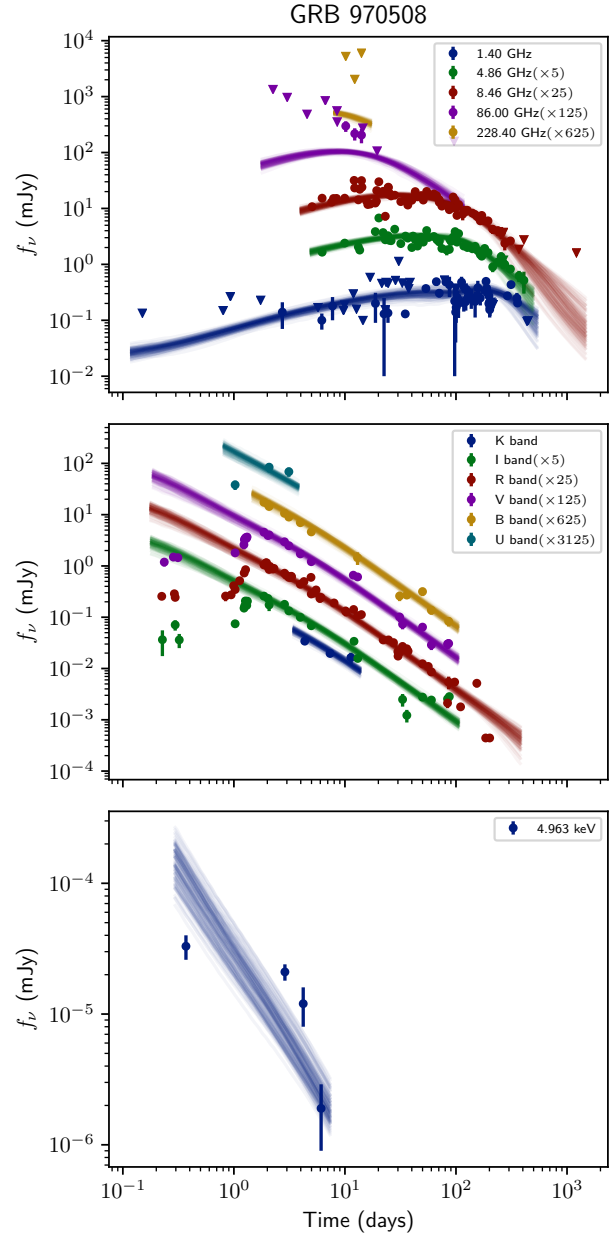


Figure A1. Fit result for GRB 970508 by sampling the GP likelihood. Observed flux density values are presented in radio, optical and X-ray bands (upper, middle and lower panel respectively) together with the posterior predictive light curves. Triangles represent $3\text{-}\sigma$ upper limits.

A sample of 100 parameter sets are randomly drawn from the inferred joint probability distribution of the parameters, and *scal-efit* light curves are drawn for each parameter set.

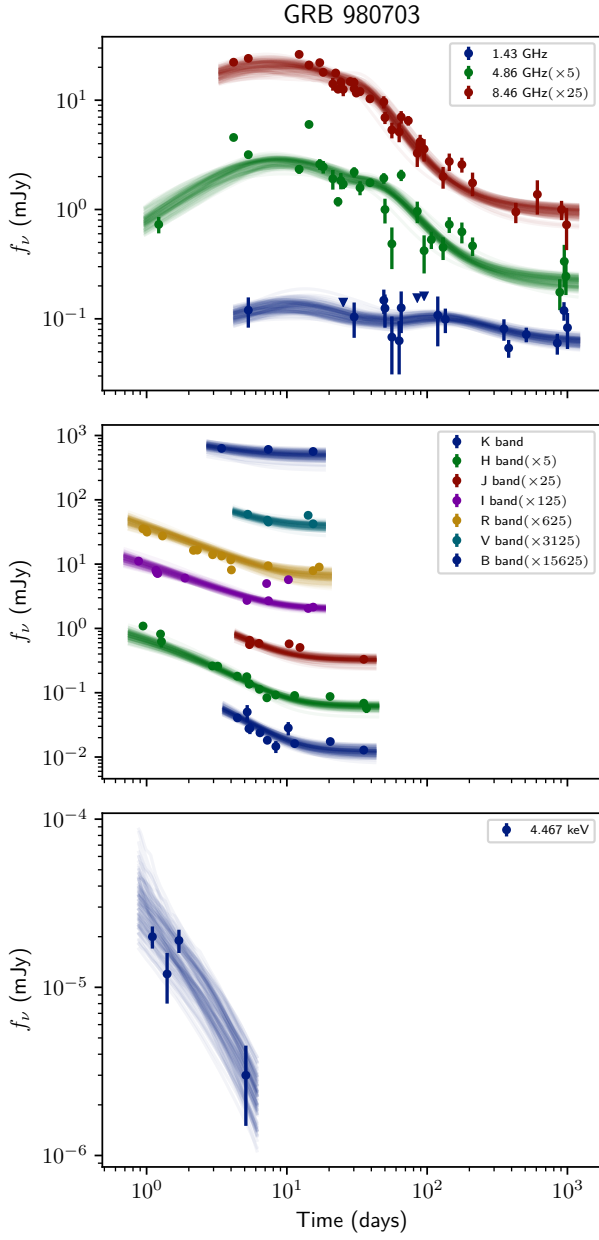


Figure A2. Fit result for GRB 980703 by sampling the GP likelihood. Observed flux density values are presented in radio, optical and X-ray bands (upper, middle and lower panel respectively) together with the posterior predictive light curves. Triangles represent 3σ upper limits.

A sample of 100 parameter sets are randomly drawn from the inferred joint probability distribution of the parameters, and `scal-efit` light curves are drawn for each parameter set. The host galaxy contribution in radio and optical is not subtracted.

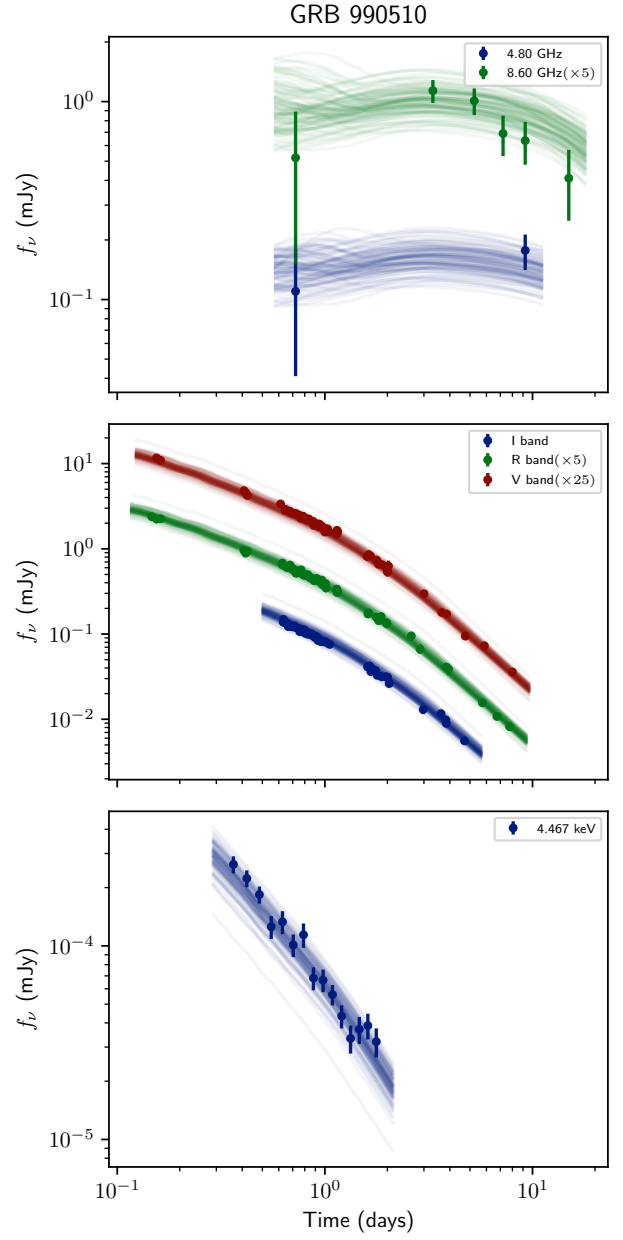


Figure A3. Fit result for GRB 990510 by sampling the GP likelihood. Observed flux density values are presented in radio, optical and X-ray bands (upper, middle and lower panel respectively) together with the posterior predictive light curves.

A sample of 100 parameter sets are randomly drawn from the inferred joint probability distribution of the parameters, and `scal-efit` light curves are drawn for each parameter set.

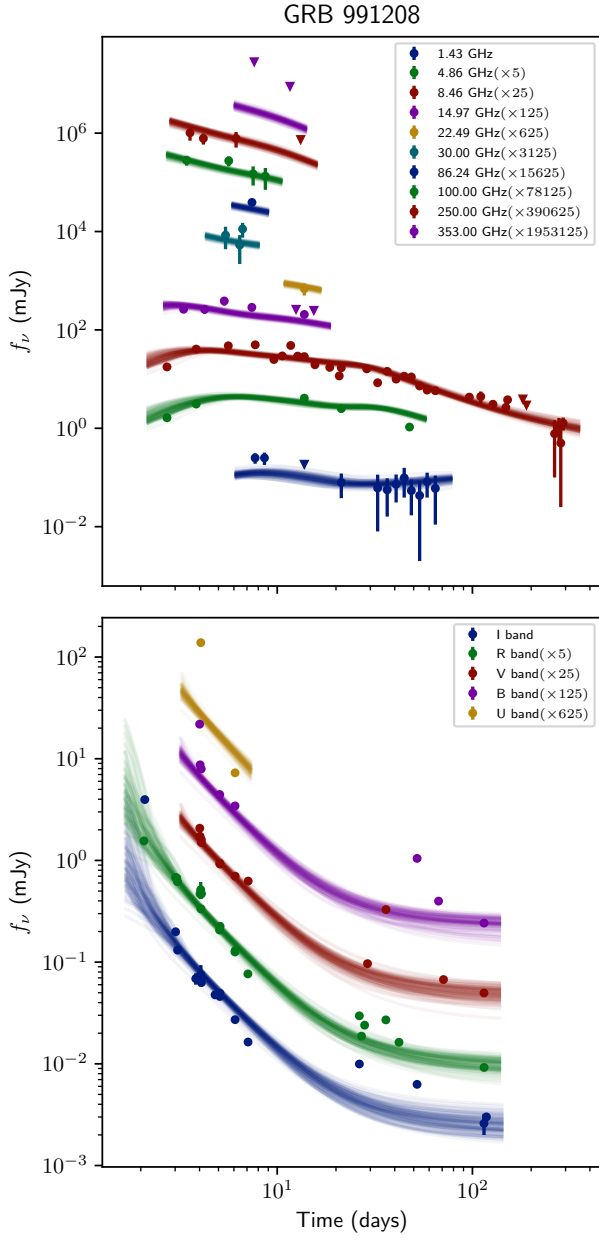


Figure A4. Fit result for GRB 991208 by sampling the GP likelihood. Observed flux density values are presented in radio and optical bands (upper and lower panel respectively) together with the posterior predictive light curves. Triangles represent $3\text{-}\sigma$ upper limits.

A sample of 100 parameter sets are randomly drawn from the inferred joint probability distribution of the parameters, and `scal-efit` light curves are drawn for each parameter set. The host galaxy contribution in optical is not subtracted.

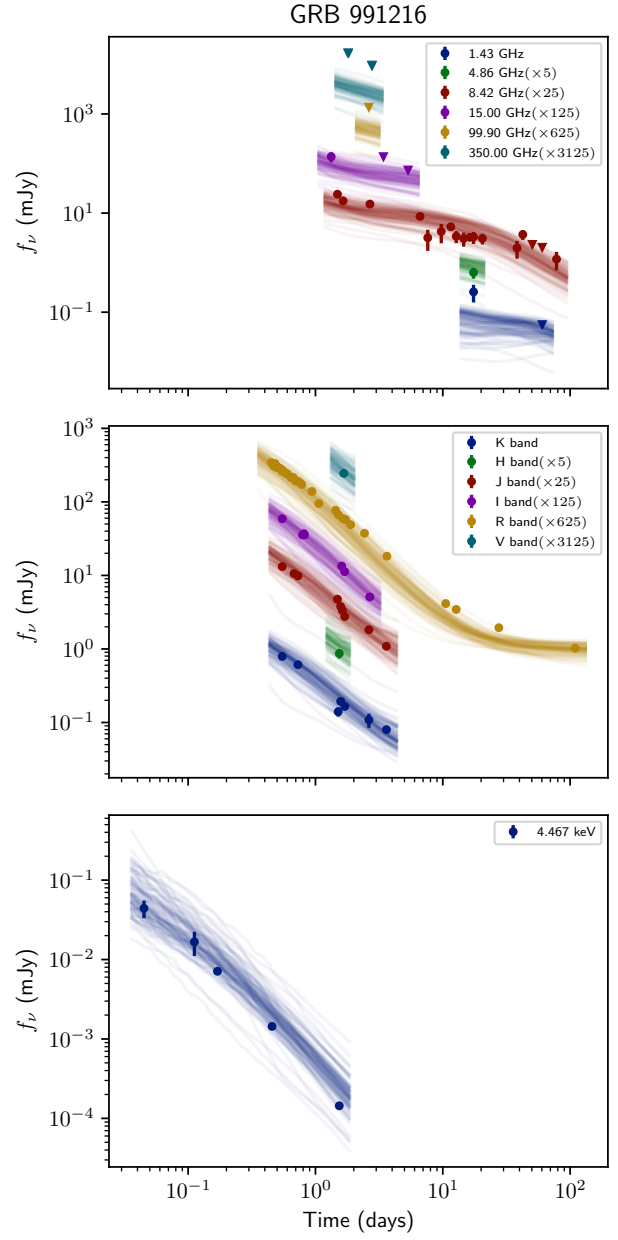


Figure A5. Fit result for GRB 991216 by sampling the GP likelihood. Observed flux density values are presented in radio, optical and X-ray bands (upper, middle and lower panel respectively) together with the posterior predictive light curves. Triangles represent $3\text{-}\sigma$ upper limits.

A sample of 100 parameter sets are randomly drawn from the inferred joint probability distribution of the parameters, and `scal-efit` light curves are drawn for each parameter set. The host galaxy contribution in optical is not subtracted.